



# Approach for the structure analysis of complex molecular systems at computer simulation

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## Abstract

The structure of complex inhomogeneous systems is a current problem in the physics of liquids, glasses, polymers, molecular biology, and material science. To understand the physical properties of such materials, one should study the structure in “micro” and “macro” levels, which demands a unified rigorous approach for structure investigations. We demonstrate that the Voronoi–Delaunay technique, which is well known in mathematics and in computer simulations, can be used for this purpose. The Voronoi–Delaunay tessellation contains complete information about the structure of a computer model. The method is applied for studying the intermediate range order to show that the behavior of the so-called prepeak in the structure factor is defined by a motif of spatial distribution of voids in the model. A large model of dense packing of spherical atoms in the process of crystallization from non-crystalline phase is analyzed. The extended linear defects of the diverse types are revealed. Investigation of the free volume distribution in the lipid bilayer in water is carried out. The results obtained can help for understanding a mechanism of diffusion of small molecules across lipid membranes.

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## 1. Introduction

Complex molecular and heterogeneous materials are objects for present day computer simulations. A typical molecular dynamic model of such systems contains tens of

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1 thousands of atoms. The models should be large enough because of large-scale struc-  
2 ture peculiarities in such materials. Different scale levels characterize the structure in  
3 this case. To study the macroscopic properties, we should take into account both the  
4 local order (nearest surrounding of atoms) and the “extended” structure. A known  
5 example of different length scale manifestation is the so-called prepeak in the struc-  
6 ture factor of some glasses and liquids obtained in diffraction experiments. This phe-  
7 nomenon was a challenge for theorists for years to explain the nature of this rather  
8 sharp peak for the completely non-crystalline materials; see for example [1] and refer-  
9 ences therein. The problem of the structure of heterogeneous materials treats a wide  
10 variety of physical and chemical properties [2]. The analysis of the structure of such  
11 systems needs a unified rigorous means of characterizing the structure in “micro” and  
12 “macro” levels. In this paper we show that the Voronoi–Delaunay technique, which  
13 is well known in mathematics and some fields of physics, can be a base for the  
14 approach [3,4].

## 15 2. The method

16 The Voronoi–Delaunay method is based on the general mathematical theorems about  
17 division of space between centers in an ensemble of discrete centers. The centers  
18 (atoms) can be distributed arbitrarily in space: orderly or disorderly, homogeneously  
19 or non-homogeneously [3,4]. Therefore, the method is claimed for studying the struc-  
20 ture of liquids and glasses, where the approaches of crystallography do not  
21 work, and for any heterogeneous material with different structure  
22 scales.

23 The main geometrical constructions of the method are the Voronoi polyhedra and  
24 Delaunay simplices. They have been used in computer simulation for structure char-  
25 acterization of the local order of atoms for many years, see e.g. Refs. [4,5]. Usual-  
26 ly they are calculated separately to get a histogram of a statistical distribution of  
27 topological or metrical characteristics of local order. However, to study the whole  
28 structure of a system, one should use the whole Voronoi–Delaunay tessellation,  
29 Fig. 1. In this case, one can analyze the spatial distribution of specific atomic config-  
30 urations to study extended structure correlations. A helpful construction for this work  
31 is the Voronoi network: the network of edges and vertices of the Voronoi polyhedra.  
32 This network permeates through the systems, and can be a background for studying  
33 structural motifs in the model. Every vertex of the Voronoi network is simultaneously  
34 the circumcenter of a Delaunay simplex. The Voronoi network lies “in the depth” of  
35 empty space between atoms, and thereby can also be used to study the free volume  
36 distribution.

37 There are different algorithms for calculation of the Voronoi polyhedra, Delaunay  
38 simplices and the Voronoi network [3,4]. Input data for this calculation is a set of  
39 coordinates of atoms for a given atomic configuration. The Voronoi network is defined  
40 by a set of coordinates of the Voronoi vertices (array {D}), and a table of connectivity  
41 of the vertices (array {DD}).

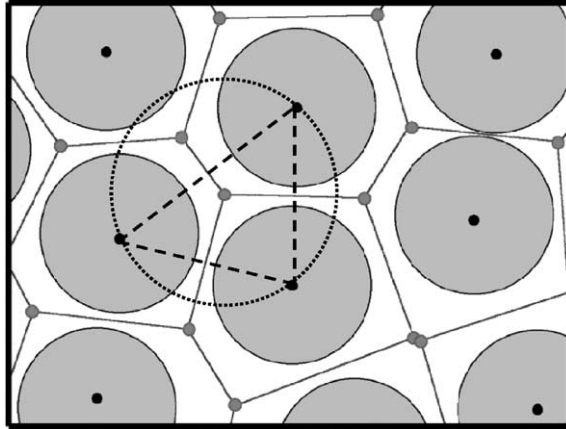


Fig. 1. Voronoi–Delaunay tessellation of a 2D atomic system: (a) *The Voronoi polyhedron* for a given atom is a volume of space, all points of which are closer to a given atom than to any other atom of the ensemble (polygons around disks). (b) *The Delaunay simplex* is a triangle (tetrahedron in 3D) defined by atoms, the circumsphere of which is empty, i.e., there are no other atomic centers inside the circumsphere (dashed lines). (c) *The Voronoi network* of a given atomic system is a network of edges and vertices of the Voronoi polyhedra (solid lines and points between atoms).

### 1 3. Investigation of the intermediate range order

3 Recently, a simple procedure to prepare models with a given intermediate order was  
 3 proposed [6]. The idea is to remove some atoms from an original packing to get a new  
 5 model with extra voids. Using different rules for removing atoms, one can get models  
 5 with different spatial distribution of extra voids, and study their physical manifestation.

7 The densest non-crystalline packing of 27,000 Lennard-Jones atoms generated by the  
 7 Monte Carlo method in a model box with periodic boundary conditions is used as an  
 9 original configuration. The atoms which are the closest to the sites of a given “stencil”  
 9 are removed. Any system of sites with a desired spatial distribution can be used as a  
 11 stencil. Fig. 2 demonstrates a model obtained after removing 15% of atoms from the  
 13 original packing. A stencil used here is a system of centers of the densest disordered  
 13 packing of spheres with radius three times larger than radius of atoms in the original  
 15 packing. So the new length scale has been added in the model. The structure factor of  
 15 the model is like the one for the original model (Lennard-Jones glass): clear main peak  
 17 and damping oscillations with increasing  $q$ . The only difference is a clear prepeak at  
 17 the small value of  $q$ , Fig. 3 (on the top). The position of the prepeak (one-third of the  
 19 position of the main peak) corresponds to the new length scale inserted in the model.  
 19 It is an interesting and non-trivial result. Indeed the extra voids in the model are very  
 21 diverse every one is a result of removing a different number of atoms (one or a few),  
 21 and the remaining atoms are not moved from their original positions. So there are no  
 23 obvious reasons to get a sharp peak in the structure factor. However, the calculation  
 23 demonstrates earnestly that the generated long-distance correlations manifest a sharp  
 peak, Fig. 3 (on the top).

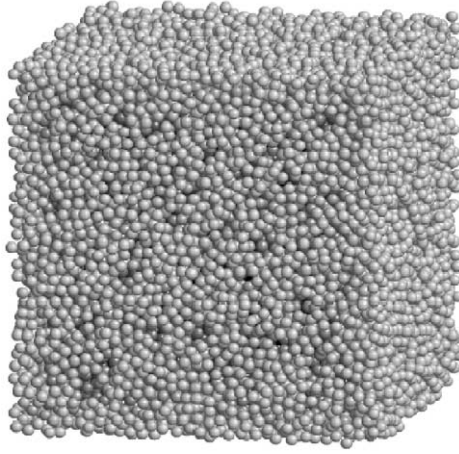


Fig. 2. Model with an intermediate order. Dense packing of spherical atoms with extra voids distributed according to a given motif, see text.

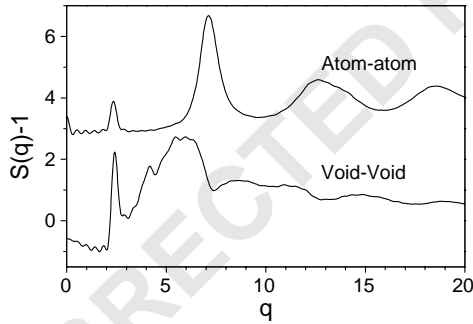


Fig. 3. Structure factors of the model shown in Fig. 2. Usual atom–atom structure factor (on the top), void–void structure factor (at the bottom), see text.

- 1 The structure factor calculated for the system  $\{D\}$  of the model is shown in
- 2 Fig. 3 (at the bottom). (It is a Fourier transformation of the radial distribution function
- 3 of the Voronoi vertices, see Fig. 1.) This function characterizes distribution on empty
- 4 space in the model, because points of  $\{D\}$  lie in the depth between the atoms. A
- 5 clear sharp peak in this void–void structure factor coinciding with the prepeak in the
- 6 atom–atom structure factor indicates that the intermediate range order is conditioned
- 7 by voids. Thus the Voronoi–Delaunay method helps us to ascertain a role of voids in
- the development of the intermediate range order.

#### 1 4. Investigation of lineal defects in model of a dense crystal

3 The other example of application of the Voronoi–Delaunay method is analysis of  
4 the model of crystal with defects. The model understudied contains 16,382 atoms in  
5 the box with periodic boundary conditions. It was prepared by molecular dynamic  
6 relaxation with Lennard-Jones potential starting from the model of hard spheres pack-  
7 ing generated in Ref. [7], where crystallization of a large disordered system of hard  
8 spheres was studied. The model obtained is a dense crystal with regions of f.c.c.  
9 and h.c.p. structure, Fig. 4. Now we demonstrate the Voronoi–Delaunay approach for  
10 analysis of the structure at the border of these regions, where a crystalline order is  
11 broken.

12 For investigation of extended structure peculiarities, one can study spatial distribution  
13 of Delaunay simplices of a specific shape. Studying defects in a crystal, it is natural to  
14 work with simplices, whose shape is atypical for a given crystalline structure. Location  
15 of such “imperfect” simplices reveals defects in the model. This approach was used  
16 in Ref. [8], where defects in a crystal after plastic deformation were studied in the  
17 computer simulation. A quantitative measure of shape can be easily defined to select  
18 the simplices, see e.g. Refs. [9,4].

19 Only a few percent of the Delaunay simplices in the model in Fig. 4 have shape,  
20 which can be marked as atypical for the given crystalline structures. We used 1.2%  
21 of the more imperfect simplices, and display their positions inside the model box,  
22 see Fig. 5. As explained in Fig. 1, the position of the Delaunay simplex is defined  
23 by the Voronoi vertex (unfortunately, these points are invisible in the scale of Fig.  
24 5). Neighboring Voronoi vertices corresponding to the imperfect simplices are con-  
25 nected by line (i.e., the Voronoi edge between such vertices is drawn). So, aggregates  
26 of the Delaunay simplices of a given type are displayed as clusters on the Voronoi  
27 network. We see these clusters of Fig. 5. The approach to show structure motifs on  
28 the Voronoi network was used for analysis of the structure of liquids and glasses in  
29 Ref. [10].

30 The clusters of imperfect simplices are very expressive, Fig. 5. Except for some  
31 complex aggregates one can see distinct clusters like straight lines, stairs, and chains.  
32 Knowing atoms in the simplices, one can obtain an arrangement of atoms correspond-  
33 ing to given clusters, Fig. 6. Line-type defect is a pile of triangles of atoms. Stairs-type  
34 defect is a pile of rhombs. Atoms in these piles are located one after another, which  
35 is unusual for the densest crystals, where atoms of one plane are located in the cav-  
36 ity between atoms of another plane. These defects are very stable and remain during  
37 a long period of relaxation of the model. The chain-like cluster is a tube of fives  
38 of atoms. It is an amusing result: we see “a five folder order ” construction in the  
39 crystal.

40 Note that we have obtained only linear extended defects. They are nuclei and inter-  
41 sections of dislocations presented in the model. The dislocations change alternation of  
42 the crystallographic planes of atoms, and define flat borders between regions of f.c.c.  
43 and h.c.p. structures, see Fig. 4. However, we do not see the flat defects in Fig. 5. It  
44 is because the Delaunay simplices in these dislocations have the typical shape for the  
45 f.c.c. and h.c.p. crystals and therefore have not been detected here.

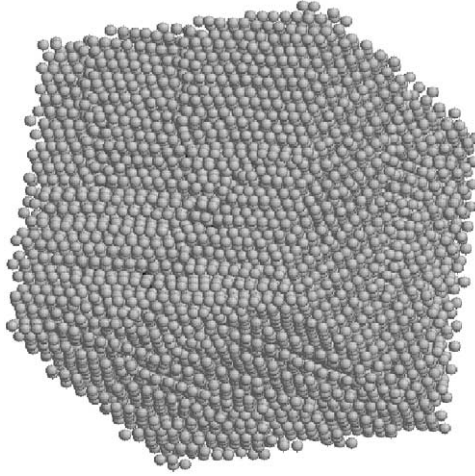


Fig. 4. Model of crystalline phase of 16,386 Lennard-Jones atoms. Breaking of interchanging crystalline planes is visible. The model contains regions of f.c.c. and h.c.p. structures.

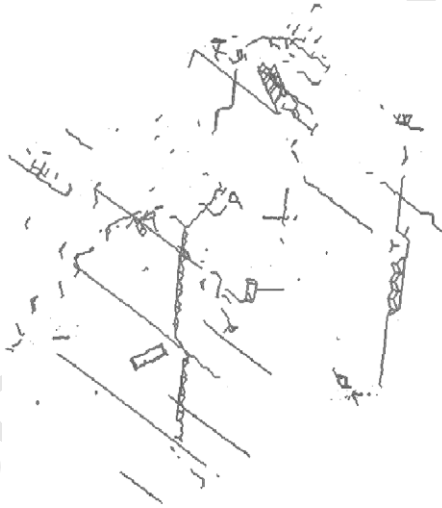


Fig. 5. Distribution of defects inside the model shown in Fig. 4. Clusters indicate the location of the Delaunay simplices, whose shape is atypical for given crystalline structures, see text.

- 1 Thus, linear defects in the molecular dynamic model of dense crystalline phase are
- 2 revealed. For this purpose the spatial distribution of the Delaunay simplices, whose
- 3 shape is atypical for the crystalline structure, was investigated.

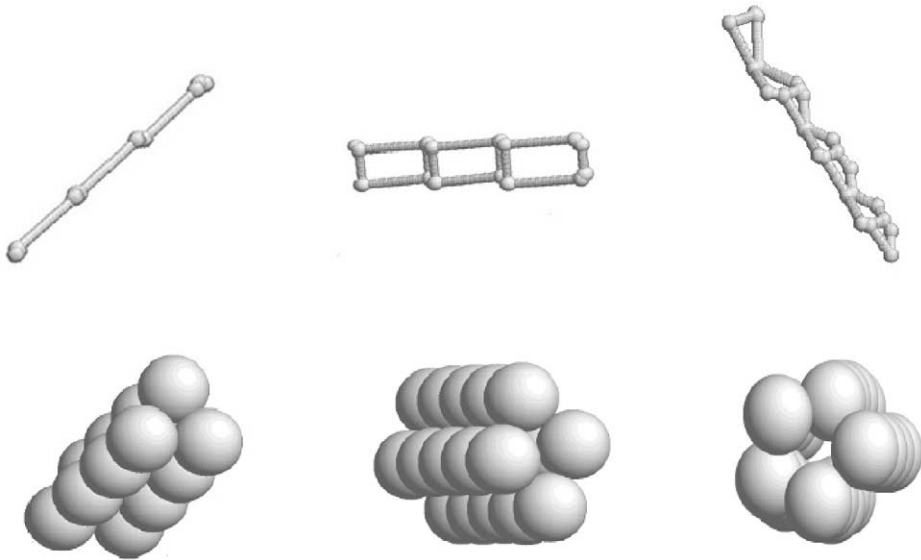


Fig. 6. Clusters on the Voronoi network for imperfect Delaunay simplices from Fig. 5 (on the top), and corresponding atomic configurations (at the bottom). Line-type (left), stairs-type (center), and chain-like defects (right).

## 1 5. Investigation of free volume distribution in a lipid bilayer

3 The molecular dynamic model of a lipid bilayer in water was generated by Schmelter  
 4 and Geiger [11]. The model box with periodic boundary conditions contains 32,125  
 5 atomic units: 200 molecules of phospholipid DPPC (dipalmitoyl-*sn*-glycero-  
 6 phosphatidylcholin) with 130 atomic units, and 6125 water molecules. The van der  
 7 Waals radii of the atomic units were used in the analysis. From a mathematical point  
 8 of view, this complex molecular system is also an ensemble of spheres. But it has some  
 9 differences in comparison with the ensembles of uniform spheres discussed above. At  
 10 first, the radii of the spheres are different: the solvent and lipid molecules consist of  
 11 different atoms. It is an important point for our problem to study free volume between  
 12 atoms. Second, some spheres of our ensemble are overlapping: the distance between  
 13 centers of chemically bonded atoms is less than the sum of their van der Waals radii.  
 14 The methods of taking all of these complexities into consideration are known and  
 15 discussed in detail in Ref. [4].

15 Distribution of the largest empty spheres inscribed between atoms gives the simplest  
 16 and obvious picture of the free volume inside the model. All empty spheres can be  
 17 found at the calculation of the Voronoi network of the model. (Remember, every  
 18 Voronoi vertex is a center of one of these inscribed spheres.) We carried out analysis  
 19 of empty spheres with radius greater than 1.4 Å. (It is a radius of water molecule, so  
 20 we look at voids where a probe-like water molecule can be placed.) It was found that  
 21 a number of such spheres are located in the mid of the bilayer: between hydrophobic

1 ends of the lipids. It is natural that the tails of the lipids do not pack densely. A  
2 number of large empty spheres are also in bulk water. It is known that water has a  
3 low dense structure with cavities of the size of the water molecule. A more unexpected  
4 result is the absence of large voids in the region of the hydrophilic ends of lipids  
5 (region of contact of bilayer with bulk water). This fact can be connected with strong  
6 interaction of atoms in this region, but it cannot exclude large voids a priori. However,  
7 the recent molecular dynamic calculation of Schmelter and Geiger demonstrates that  
8 water molecules penetrate into this region. Seemingly, these molecules fill in all large  
9 voids in this area [12]. All this information is helpful for understanding the mechanism  
10 of diffusion of small molecules across lipid membranes.

## 11 6. Conclusion

12 The Voronoi–Delaunay method is an extremely helpful tool for analysis of the struc-  
13 ture of computer models of complex molecular systems. It enables one to study both  
14 local order of atoms and extended structure correlations of atoms, as well as voids  
15 between atoms.

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## 17 References

- 18 [1] S.R. Elliott, *J. Phys.: Condens. Matter* 4 (1992) 7661 (Printed in the UK).
- 19 [2] S. Torquato, *Random Heterogeneous Materials: Microstructure and Macroscopic Properties*, Springer, Berlin, 2002, in Press.
- 20 [3] A. Okabe, B. Boots, K. Sugihara, S.N. Chin, *Spatial Tessellations: Concepts and Applications of Voronoi Diagrams*, Wiley, Chichester, 2000.
- 21 [4] N.N. Medvedev, *Voronoi–Delaunay Method for Non-Crystalline Structures*, SB Russian Academy of Science, Novosibirsk, 2000 (in Russian).
- 22 [5] J. Finney, *Roy. Soc. London* 319 (1970) 479.
- 23 [6] V.P. Voloshin, S. Beaufils, N.N. Medvedev, *J. Mol. Liquids*, 2002, in Press.
- 24 [7] V.A. Luchnikov, A. Gervois, P. Richard, L. Oger, J.-P. Troadec, *J. Mol. Liquids*, 2001, in Press.
- 25 [8] I.I. Gainutdinov, Yu.T. Pavlyukhin, *Doklady RAS*. 364 (2) (1999) 203 (in Russian).
- 26 [9] V.A. Luchnikov, N.N. Medvedev, Yu.I. Naberukhin, V.N. Novikov, *Phys. Rev. B* 51 (1995) 15,569.
- 27 [10] I.Yu. Naberukhin, V.P. Voloshin, N.N. Medvedev, *Mol. Phys.* 73 (4) (1991) 917.
- 28 [11] R. Schmelter, A. Geiger, in preparation.
- 29 [12] A. Geiger, *Horizons of Complex Systems* (4-8 December), Messina, Italy, 2001.